

1 **The impact of high Zn⁰ concentrations on the application of**
2 **AGNES to determine free Zn(II) activity**

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8 **Abstract**

9 AGNES (Absence of Gradients and Nernstian Equilibrium Stripping) determination of
10 free Zn(II) in a solution can be affected by the reaching of high Zn⁰ concentrations
11 inside the amalgam. At concentrations about the solubility limit of Zn⁰ in mercury, the
12 formation of dendrites and powders around the mercury surface can be seen with an
13 optical microscope. At concentrations of Zn⁰ quite below the solubility limit, an
14 anomalous stripping current appears which increases with decreasing supporting
15 electrolyte concentration. The current along the stripping time exhibits a convex shape,
16 which is labelled here as “anomalous convex behaviour” (acb). The origin of acb is
17 tentatively ascribed to different kinetic reasons (amongst which the electroneutrality
18 limitation due to low ionic strength outstands), but more than one cause is necessary for
19 a full account of the experimental observations. With various strategies, like monitoring
20 the charge as response function or by application of lower gains, AGNES can
21 successfully probe these high concentrations.

22
23 **Keywords:** AGNES, stripping analysis, amalgam, free metal.

24

25 **1. Introduction**

26 AGNES (Absence of Gradients and Nernstian Equilibrium Stripping) is an emerging
27 electroanalytical technique specifically designed to determine free metal concentrations
28 [1]. Its application to a variety of systems, ranging from synthetic solutions to
29 Mediterranean seawater, has been extensively validated on theoretical grounds or with
30 consolidated techniques such as Ion Selective Electrode, Resin Titration or Scanned
31 Stripping Chronopotentiometry [2-9]. The key idea of AGNES is the preconcentration
32 of reduced metal (Zn⁰ in this work) inside the amalgam up to the equilibrium value (set
33 by the applied potential and Nernst equation). The determination of the amount of Zn⁰
34 at equilibrium has been performed, up to now, via the stripping current under diffusion
35 limited conditions. At large concentrations of the reduced metal (and low ionic
36 strength), we have observed a current distortion, which might hinder the standard
37 application of AGNES in some systems (for instance, in determining the complexation
38 capacity of wine with the standard HMDE [4] or when working with film electrodes
39 [9]).

40

41 A distortion of the current is clearly expected when the solubility limit of Zn⁰ inside the
42 amalgam is exceeded (5.8% in atomic percentage [10], approximately 4 M, at 25°C). It
43 is obvious that stripping currents for aimed Zn⁰ concentrations larger than the solubility
44 limit could not be used for the analytical general purpose of AGNES. But, a literature
45 search indicates that problems might appear well before. For instance, Arevalo et al.
46 [11] found a linear behaviour of the stripping current with [Zn⁰] at least up to
47 [Zn⁰]=3×10⁻³ M, while reporting that Babkin had found that the linear relationship
48 ceased from [Zn⁰]=1.2×10⁻² M onwards (the current increased more rapidly than [Zn⁰]).
49 Moreover, a long tradition of literature [12-15] has pointed out that there might be
50 “kinetic” problems: even at reduced metal concentration globally lower than the

51 solubility limit, during the deposition there might be spots (in the amalgam, close to the
52 surface of the interphase) where the local concentration exceeds the solubility limit and
53 forms deposits. More recently, experimental and theoretical evidences (with Tl) indicate
54 that the stripping currents are strongly influenced by low supporting electrolyte
55 concentrations [16,17].

56

57 The outline of this article is as follows. We summarize AGNES methodology and
58 introduce two new concepts: i) the time-function η , expressing the direct proportionality
59 between stripping current and reduced metal concentration, which allows for diagnosing
60 anomalous behaviours and ii) the charge as response function in AGNES. We describe
61 the case of formation of Zn⁰ deposits around the mercury drop linking optical and
62 electrochemical information. We study the deviation from the current-concentration
63 linearity, consider various hypotheses which might participate in the justification of the
64 experimental results and describe practical strategies to apply AGNES when dealing
65 with high concentrations.

66

67 **2. The principles of AGNES**

68 AGNES is a stripping technique whose novel idea is the attainment of a special
69 situation of equilibrium by the end of the deposition stage [1]. The sought equilibrium
70 can be achieved by application of a deposition potential E_1 for a sufficiently long time t_1
71 (see table 1). E_1 is usually just a few millivolts more negative than the standard formal
72 potential of the couple $E^{0'}$, so that the prescribed gain Y , attained by the end of the
73 deposition stage and given by Nernst law,

$$74 \quad Y = \frac{[\text{Zn}^0]}{[\text{Zn}^{2+}]} = \exp\left[-\frac{2F}{RT}(E_1 - E^{0'})\right] \quad (1)$$

75 is moderate. F is the Faraday, R the gas constant, T the temperature and the
76 concentrations refer to the homogeneous profile (i.e. absence of gradients) inside the
77 mercury electrode and in the solution. $[Zn^0]$ in this article refers exclusively to the bulk
78 homogeneous concentration obtained inside the drop by the end of the first stage, so that
79 it is not a time-dependent variable.

80

81 Up to this work, the only response function of AGNES considered has been the
82 intensity current at a fixed time t_2 (within the stripping stage) which is sought to be
83 under diffusion limited conditions (with a stripping gain Y_2 computed with eqn (1), but
84 replacing E_1 with E_2 , a much more positive potential used for the reoxidation stage).
85 These conditions, due to the linearity of the continuity equation for Zn^0 inside the
86 amalgam (see appendix in [1]), ensure that the faradaic current I (obtained from the
87 measured current once a suitable blank has been subtracted) is proportional to $[Zn^0]$:

$$88 \quad I = \eta [Zn^0] \quad (2)$$

89 where the new proportionality factor η only depends on the elapsed time of the stripping
90 step (which –for simplicity- we also label t_2) and on the characteristics of the diffusion
91 of Zn^0 in the amalgam (diffusion coefficient and shape and size of the drop). The plot of

$$92 \quad \eta(t_2) = \frac{I(t_2)}{[Zn^0]} = \frac{I(t_2)}{Y[Zn^{2+}]} \quad (3)$$

93 along the stripping time should be the same –for a fixed electrode and metal- regardless
94 of Y and $[Zn^{2+}]$. In order to control deviations in the stripping current from the expected
95 linear behaviour, we will plot $\eta(t_2)$, to which we will refer here as “plot of normalised
96 currents”. The “reference curve” will be given by the collapse of all plots of $\eta(t_2)$ at
97 conditions with no anomalies (usually low metal concentrations and high ionic
98 strength).

99

100 Combining eqns. (1) and (2), we can write

$$101 \quad I = \eta [Zn^0] = \eta Y [Zn^{2+}] = h [Zn^{2+}] \quad (4)$$

102 where $h(t_2)$ is the proportionality factor between the experimental faradaic current and
103 the sought free metal concentration in the solution bulk. For analytical purposes, h can
104 be determined from a calibration plot and, then, used with the faradaic current measured
105 in the studied sample.

106

107 Alternatively to the use of the intensity current (at a fixed t_2), we explore here the use of
108 the deposited charge. Indeed, by simple integration of eqn. (2) along the (complete)
109 stripping time:

$$110 \quad Q = \int_0^\infty I dt_2 = \int_0^\infty \eta [Zn^0] dt_2 = \left(\int_0^\infty \eta dt_2 \right) [Zn^0] = \eta_Q [Zn^0] \quad (5)$$

111 where we have added a subscript Q to indicate that the proportionality refers to the
112 charge (and not to the current, as usual). η_Q can be related to the volume of the
113 electrode. If we assume a spherical electrode of radius r_0 :

$$114 \quad \eta_Q = 2F \frac{4}{3} \pi r_0^3 \quad (6)$$

115

116 Combining eqns. (1) and (5), we can write

$$117 \quad Q = \eta_Q [Zn^0] = \eta_Q Y [Zn^{2+}] = h_Q [Zn^{2+}] \quad (7)$$

118

119 If one has determined η_Q from a calibration plot, further charge measurements in a
120 sample solution can yield its free metal concentration:

$$121 \quad [Zn^{2+}] = \frac{Q}{Y\eta_Q} \quad (8)$$

122 In principle, one could use any measure of the faradaic charge (i.e. in the first or second
123 stage), but we have noticed that the stripping charge is much more reproducible than the
124 deposition charge, and, so, in this work, Q always refers to the charge measured in the
125 stripping stage.

126

127 **3. Materials and Methods**

128 **3.1 Reagents**

129 Zinc and cadmium stock solutions were prepared from $\text{Zn}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}$ and
130 $\text{Cd}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}$ (Merck, analytical grade), respectively, and standardised by means of a
131 complexometric endpoint titration with EDTA [18]. Diluted solution of Zn and Cd were
132 prepared from the Zn and Cd stock solutions, respectively, or from Zn and Cd 1000
133 mg/L standard solution (Merck). Potassium nitrate was used as inert supporting
134 electrolyte and prepared from solid KNO_3 (Fluka, TraceSelect). The ionic strength has
135 been mainly fixed at 0.05 M (because of the interest of wine matrices [4,19,20]. Titrisol
136 (Merck) or standard 0.1 M solutions (Riedel de Haen) of KOH and HNO_3 were added to
137 fix the pH at the desired values. Mercury (Fluka p.a.) was used in the working electrode.

138

139 Ultrapure water (Milli-Q plus 185 System, Millipore) was employed in all the
140 experiments. Purified water-saturated nitrogen $\text{N}_2(50)$ was used for deaeration and
141 blanketing of solutions.

142

143 **3.2 Instrumentation**

144 Voltammetric measurements were carried out with Eco Chemie Autolab PGSTAT30
145 PGSTAT12 or PGSTAT10 potentiostats attached to a Metrohm 663 VA Stand and to a
146 computer by means of the GPES 4.9 (Eco Chemie) software package. The working

147 electrode was a Metrohm multimode mercury drop electrode. The smallest drop in our
148 stand has been chosen, which, according to the catalogue, corresponds to a radius
149 around $r_0 = 1.41 \times 10^{-4}$ m. This size was confirmed by measurements with our
150 microscope. The auxiliary electrode was a glassy carbon electrode and the reference
151 electrode was Ag | AgCl | (3 mol L⁻¹) KCl, encased in a 0.1 mol L⁻¹ KNO₃ jacket.

152

153 A glass combined electrode (Orion 9103) was attached to an Orion Research 720A
154 Ionanalyzer and introduced in the cell to control the pH (all the experiments were
155 performed at pH in the range 3-5). Glass jacketed cells (thermostated at 25.0°C)
156 provided by Metrohm or Afora were used.

157

158 To monitor the evolution of the mercury drop, a CCD digital camera (Lumenera
159 INFINITY2-1) was coupled to a microscope Navitar Zoom 6000 (with coaxial
160 illumination), following previously reported designs [21]. To obtain a large
161 magnification range (6.96x – 228.6x), two different objective lenses (10x and 50x,
162 working distances 33 mm and 13 mm respectively) were used. A linear motor
163 positioning stage (Standa, Lithuania) controlled by a microstep stepping motor driver
164 was used to move the microscope in 3D with a resolution of ca. 0.16 μm. A flat quartz
165 window was added to a double-wall voltammetric cell.

166

167 3.3 Procedures

168 AGNES experiments were applied with the simplest potential program, i.e. a fixed
169 constant potential during the deposition step (see [1] for more details). The standard
170 response function of AGNES has been the value of the current at a certain time t_2 in the
171 reoxidation step (e.g. 200 ms for these experiments with this HMDE). As an alternative

172 source of information, we compute the charge along the reoxidation step by just
173 integrating the current measured along the second stage using the trapezoid rule (and a
174 zero current for $t_2=0$, which will imply an underestimation of the actual charge):

$$175 \quad Q = \int_0^{\infty} I dt_2 \approx \Delta t \left(\sum_{j=1}^{n_{\max}} I_j \right) \quad (9)$$

176 where Δt is the sampling time (either 1 ms or 50 ms), I_j is the faradaic current at a given
177 measurement point j corresponding to the variable time t_2 (i.e. $j= t_2/\Delta t$). The faradaic
178 current has been obtained by subtracting a blank current (usually negligible in these
179 experiments) and I_{∞} (which is the residual current, average of the measured current
180 around the end of the second stage). Typically, the stripping stage lasted 50 s, so that the
181 standard maximum number of points is $n_{\max}= 50/0.050=1000$. Due to the discretization,
182 the actual value of Q –and the ensuing parameters h_Q and η_Q – depends on the
183 experimental interval time Δt , converging to a fixed value for low enough Δt .

184

185 In order to estimate an upper bound for the maximum possible charge in a stripping
186 stage with standard behaviour, we have assumed cottrellian behaviour up to the first
187 sampled point (whose current is labelled here as I_1), given that sphericity and finite
188 volume effects can be considered negligible for that time. Then, the maximum charge
189 up to the first sampled point should be

$$190 \quad Q_1 = \int_0^{\Delta t} I dt = \int_0^{\Delta t} 2F4\pi r_0^2 [Zn^0] \sqrt{\frac{D_{Zn^0}}{\pi t}} dt = 2 \times 2F4\pi r_0^2 [Zn^0] \sqrt{\frac{D_{Zn^0} \Delta t}{\pi}} = 2 I_1 \Delta t \quad (10)$$

191 where D_{Zn^0} stands for the diffusion coefficient of Zn^0 . Under this overestimating
192 assumption, the charge would contain an extra term:

$$193 \quad Q_{\text{bound}} \approx \Delta t \left(\frac{3}{2} I_1 + \sum_{j=1}^{n_{\max}} I_j \right) \quad (11)$$

194

195 Another voltammetric technique used here, the Linear Anodic Stripping Voltammetry
196 (L-ASV) [13], consists in the application of a first deposition step at a constant potential
197 followed by a linear scan of the potentials at a scan rate of 17 mV/s, moving from the
198 deposition to the reoxidation potential.

199

200 **4. Results and discussion**

201 **4.1 Observations**

202 In order to proceed to a systematic discussion of the results, we summarize some
203 regularities (despite increasing irreproducibility with higher concentrations) of a large
204 number of experiments about the application of the simplest program of AGNES (see
205 table 1) to concentrated solutions of Zn(II) in 7 observations:

206

207 • Observation 1: *For high concentrations, I is no longer proportional to [Zn²⁺] or*
208 *[Zn⁰]. The calibrations of I vs [Zn⁰] (or [Zn²⁺]) follow the linear behaviour expected*
209 *and reported in previous works [2-4,6], but the linearity ceases at high concentrations.*
210 *See, for instance, in Fig 1, how the calibration plot for the currents with [KNO₃] =0.05*
211 *M (full triangle markers) bends at concentrations above [Zn⁰] ≈ 4×10⁻⁴ M (notice that*
212 *[Zn⁰] is a key parameter for the stripping stage). Given that the gain was 50, this*
213 *corresponds to [Zn²⁺]≈8×10⁻⁶ M. The typical stripping current, predicted by the eqns.*
214 *in appendix of ref. [1] and experimentally seen in previous AGNES literature, follows a*
215 *fast decay, as the one shown by the reference line in the normalized currents plot (thick*
216 *lines in Figs 2 and 3 or circle markers in Fig 4). In the anomalous stripping at high*
217 *concentrations (see, for instance, Figs 2, 3, 4 and 5), there is a convex region of the*
218 *stripping current (I vs t₂), which, for brevity we label here as “anomalous convex*
219 *behaviour” (acb). Experiments with larger amounts of supporting electrolyte (e.g. 0.5*

220 M) exhibit a less pronounced acb effect: see how in Fig 1, the linearity of the current -as
221 commented above- is kept until $[Zn^0] \approx 4 \times 10^{-4}$ M for $[KNO_3] = 0.05$ M (full triangle
222 markers), while for $[KNO_3] = 0.5$ M (full circle markers) the linearity is practically kept
223 until $[Zn^0] \approx 0.05$ M (although for this $[Zn^0]$ the acb is observed, the current at $t_2 = 200$
224 ms approaches the expected value). The restoration of the standard stripping behaviour
225 with higher ionic strength can also be seen in Fig 3.

226

227 • Observation 2: *Q* is linear with $[Zn^0]$ for a much larger range than the intensity
228 current measured at a given t_2 . The calibrations using charge (instead of the intensity
229 current) as the response function of AGNES keep the linearity even at the
230 concentrations for which *I*-plots bend (compare full and open markers in Fig 1).
231 Linearity of the charge is not yet lost around $[Zn^0] \approx 5 \times 10^{-3}$ M for $[KNO_3] = 0.05$ M or at
232 $[Zn^0] \approx 0.2$ M for $[KNO_3] = 0.5$ M. Obviously, the charge would no longer be useful when
233 Zn^0 saturation is reached inside the drop. However, there is loss of linearity even before
234 the computed saturation limit.

235

236 • Observation 3: *Sluggish stripping*. For huge concentrations, larger than those
237 probed in Fig 1 (e.g. $[Zn^{2+}] = 0.016$ M and $Y = 30$ leading to $[Zn^0] = 0.48$ M; see Fig 5), the
238 initial recorded current is much less than expected (in Fig 4 see how the normalized
239 currents denoted with diamond markers are below the reference curves depicted with
240 circle markers). The current is initially almost “levelled-off” to a practically constant
241 value for several seconds (see Fig 5) and, so, we call “sluggish stripping” to this acb
242 variant or feature. This levelling-off effect appears (at high concentrations) when *Q* is
243 still linear with $[Zn^0]$ and remains for larger concentrations. The levelling-off effect
244 tends to decrease with increasing ionic strength.

245 • Observation 4: *Initial overcoming of the diffusion limited flux* (enhancement of
246 the current or exaltation [22]). For intermediate-high concentrations, the anomalous
247 stripping current is –at least at some points with short t_2 - higher than the diffusion
248 limited expected one (see upper curves in Figs 2, 3 and 4). Later in the stripping process
249 (i.e. for longer t_2) the normalized representation (i.e. the experimental $\eta(t_2)$) becomes
250 lower than the reference current (so that, from $t_2=0$ to $t_2=50$ s, the same charge is
251 stripped off in all cases within the range where Q is linear with the concentration) as
252 seen, for instance, in the diamond series in Fig 2 or 3. In order to experimentally
253 confirm that the enhanced current of the first measured points is not the compensation
254 of a very small initial current (at shorter stripping times than the first probed point), we
255 have repeated some experiments with a shorter interval time ($\Delta t=1$ ms instead of the
256 typical 50 ms value), and, we could see (Fig 4) that the enhanced currents are in
257 operation from 1 ms onwards in this case. To confirm that the enhanced currents
258 actually correspond to a stripping process faster (at some time of the stripping) than that
259 of diffusion limited conditions, we have plotted in Fig 6 the normalized stripped charge
260 $Q(t_2)/[Zn^0]=\eta_Q(t_2)$ in front of the stripping time: it is clearly seen that, say at $t_2=0.5$ s,
261 the stripped charge of the series for $Y=2$ (square markers) even overcomes the upper
262 bound computed with eqn. (11) (diamond markers).

263

264 The enhancement effect disappears for sufficiently high ionic strength and increases
265 with a more positive stripping potential: see how in Fig 3 the exaltation is larger for
266 $Y_2=10^{-12}$ than for $Y_2=10^{-8}$. For Cd, we have not found current enhancement with respect
267 to diffusion limited conditions.

268

269

270 • Observation 5: *Longer deposition times can remediate the enhanced currents.*

271 The anomalous convex behaviour with enhanced current, for intermediate-high
272 concentrations of amalgamated metal, disappears at long deposition times t_1 as seen in
273 Fig 2 (following the curves downwards, the deposition time increases and the currents
274 eventually collapse with the thick reference line representing diffusion limited
275 conditions).

276

277 • Observation 6: *Longer deposition times cannot remediate the sluggish stripping.*

278 When dealing with huge [Zn⁰] (leading to observation 3), we have not found a trend of
279 the recorded currents towards the reference curve when increasing the deposition time
280 (see Fig 5).

281

282 • Observation 7: *Formation of powders and dendrites.* With the microscope, we
283 have observed them around the surface of the mercury electrode using a deposition
284 potential in the limit of diffusion all along the first stage of AGNES ($Y = 10^{10}$) leading
285 to conditions close to saturation inside the amalgams. Calusaru and Kuta [23-25] have
286 previously reported this kind of phenomenon for metals like Cu, Au and Cd in mercury.
287 But, to our knowledge, such phenomenon has not yet been described for Zn. In general,
288 we have observed that this phenomenon of deposition in a kind of external shell (or
289 shield) is quite irreproducible, not only in the morphology of the dendrites, but also in
290 the moment that it appears.

291

292 At ionic strength 0.05 M, the solid deposition can be detected visually at concentrations
293 around $[Zn^{2+}] = 2 \times 10^{-2} M$. First, some bubbles appear, while at longer deposition times
294 the formation of crystalline structures (dendrites) starts (see Fig 7). These dendrites

295 dissolve in the drop in some cases, while in other instances they grow becoming bigger
296 and moving to the bottom of the drop, and leading occasionally to the falling of the
297 drop. In the cases where we find solid deposition, the second stage shows the sluggish
298 stripping phenomena.

299

300 In order to discard irreversibility of Zn as participant in the solid deposition on the Hg
301 drops, we checked its appearance with Cd. We observed that the phenomenon of
302 precipitation was more reproducible with Cd and that it happened at higher
303 concentrations than for Zn, at [Cd²⁺]=0.2 M. For Cd (data not shown), the process starts
304 with the formation of a powder around the drop, followed by a movement which looks
305 like a “boiling surface”, where we can see the formation of some bubbles. At longer
306 times a “shell” of dendrites is formed, with the ability to move around the drop. This
307 can produce the loss of contact between the capillary and the drop.

308

309 4.2 Tentative interpretations

310 Shell formation of powders and dendrites (observation 7) was extensively explained in
311 terms of a quantum effect of deposit without mechanical contact [23], but –as saturation
312 is likely to be attained in our case- a simpler explanation is the solid deposit on
313 precipitated Zn⁰ [24,25]. The 6 first observations of previous sub-section could arise
314 from just one cause or from a combination of causes. We next discuss whether the 6
315 observations can be explained with just one cause, given that some observations are
316 clearly connected (e.g. when acb appears due to either observation 3 or 4, there will be
317 the loss of linearity described in observation 1) or some observations can be thought as
318 variants of a general process (e.g. the enhanced stripping currents progressively tend –
319 for higher Zn⁰ concentrations- towards the sluggish stripping behaviour, as seen in Fig

320 4, without a clear discontinuity in the morphology of the stripping curves). The list of
321 considered hypothetical causes is:

322 i) Turbulence hypothesis. Convective flows in the amalgam can explain the intensity
323 current exaltation.

324 ii) Electroneutrality limitation. The low ratio between the supporting electrolyte
325 concentration and the active species concentration can result in a conductivity
326 limitation.

327 iii) Irreversibility of Zn⁰ reoxidation. This might explain the sluggish stripping.

328 iv) The “island” hypothesis. The formation of a biphasic amalgam (with crystals close
329 to the surface) can produce exaltation or sluggish stripping.

330 v) The shell hypothesis. The formation of a shell around the drop can explain the
331 sluggish stripping.

332 vi) Intermetallic formation. Impurities of Cu present in the mercury or in the solution
333 could result in the formation of intermetallic compounds.

334

335 Given that hypotheses iii) to vi) are less plausible, they are discussed in the Appendix,
336 while we focus here on the first two hypotheses.

337

338 i) The “turbulence” hypothesis [22,26-29]. Observation 4 requires some mechanism to
339 produce currents larger than those yielded by diffusion limited conditions. Given that
340 migration of Zn⁰ has to be discarded because of its null charge, one should resort to
341 convection. One might think, then, of a streaming flow inside the drop during the
342 stripping stage. The driving force for this internal flow could be the dramatic change in
343 density during the reoxidation of large concentrations of Zn⁰ (perhaps interacting with
344 the density gradient between the amalgams in the drop and in the capillary) or changes

345 in the surface tension. We could accept that under stagnant uniform conditions, the
346 reoxidation might proceed without turbulence (despite the density gradient or other
347 driving force), but any disturbance (let it be denoted with the generic name “seeds”)
348 might trigger the appearance of convective flows. These seeds could be remaining flows
349 generated along the deposition step, remaining irregularities such as very small islands,
350 patches of different surface tension, viscosity or potential or any other inhomogeneity.
351 Once a small turbulence is initiated by a seed, it might suffer a kind of autocatalytic
352 growth. Within this interpretative framework, if t_1 increases, the seeds disappear (more
353 homogeneization) and the anomalous behaviour disappears (observation 5). Indeed, for
354 larger deposition times, the streaming behaviour -in the stripping stage- would start
355 later, due to the progressive extinction of seeds of streaming remaining from the
356 deposition stage. The additional peak in Linear ASV experiments (see Fig 8 and the
357 island hypothesis in the appendix for more details) could be due to such a convective
358 transport.

359

360 Kolthoff and coworkers [22], reported that the streaming maxima tend to disappear at
361 larger ionic strengths and interpreted this observation following von Stackelberg and
362 Doppelfeld as the potential distribution being affected by the background electrolyte
363 [27]. The screening effect of the capillary can result in an inhomogeneous electric field
364 around the drop which would lead to the exaltation [30]. More recently, Islam *et al.* [26]
365 reported visual confirmation of the streaming process in a HMDE. The lack of enhanced
366 currents in Cd could be linked to the shorter distance between the stripping potential
367 and the potential of zero charge [22,27] for Cd in comparison with Zn.

368

369 However, this hypothesis of internal turbulence cannot explain -without an additional
370 phenomenon- observation 3 (levelling-off effect), because any amalgam turbulence
371 always enhances the (initial) current.

372

373 ii) The electroneutrality limitation. When the potential is suddenly stepped from the
374 deposition potential to the stripping potential, a huge amount of positively charged Zn²⁺
375 ions is delivered into the solution. One might think of a limitation in the supply of
376 anions from the bulk of the solution to compensate for the released cations, thus
377 producing an effective potential difference across the electrodic interfase which is quite
378 different from the one applied, so that diffusion limited conditions inside the drop
379 would not be achieved. This could explain the sluggish stripping seen at huge
380 concentrations and its practical constancy during a relatively long period (observation
381 3). Further support comes from the observation that larger concentrations of supporting
382 electrolyte favour (see Figs 1) the keeping of the linearity of the current, as commented
383 above around observation 1 and the decreasing of the sluggish stripping (see Figs 3 and
384). The similarity of the sluggish stripping of Zn and Cd at comparable high
385 concentrations also lends support to the origin of acb in electroneutrality. This
386 explanation for the sluggish stripping (observation 3) matches theoretical and
387 experimental results (on Tl) recently published [16,17] and could also imply that the
388 potential relevant for the Nernstian equilibrium be different from the nominal applied
389 potential. However, this mechanism of low medium conductivity cannot explain (at
390 least straightforwardly) the current enhancement of observation 4 (for Zn) or why
391 longer deposition times reduce them (observation 5) given that, for the stripping stage,
392 we impose diffusion limited conditions inside the amalgam where migration cannot
393 have any effect on Zn⁰ transport. In experiments with different stripping potentials, we

394 observed a larger current enhancement (with respect to the reference curve which
395 corresponds to diffusion limited conditions inside the amalgam) for $Y_2=1\times 10^{-12}$ than for
396 $Y_2=1\times 10^{-8}$ (see Fig 3). We speculate that with $Y_2=1\times 10^{-12}$ a larger turbulence in the
397 amalgam is produced, perhaps due to a migration turbulence in solution transmitted by
398 mechanical contact to the amalgam, i.e. migration of Zn^{2+} away from the electrode
399 would *induce* convection inside the amalgam, this also explaining why all acb effects
400 disappear at sufficiently high ionic strength.

401

402

403 From the above discussion, it seems that a combination of causes is required to explain
404 all the observations. A medium conductivity limitation (hypothesis ii) seems key, given
405 the dramatic supression of the acb effect for sufficiently high ionic strength and a very
406 reasonable explanation of observations 1 and 2. Moreover, the limitation of
407 conductivity easily accounts for the sluggish stripping of observation 3 or levelling-off
408 effect and observation 6. However, the currents higher than those of diffusion limited
409 conditions in observation 4 require some enhancement mechanism, which is not
410 inherent to an electroneutrality limitation. So, we could accept the turbulence
411 hypothesis (i) as adequate and also likely to be involved in observation 5 as explained
412 before: some homogeneization of the system with increasing deposition time might tend
413 to deactivate the seeds of the turbulence (in the generic sense that includes potential or
414 density inhomogeneities, islands, etc.). The lack of enhanced currents for sufficiently
415 high ionic strength suggests the idea of the turbulences being induced (or maintained)
416 by the migration of Zn^{2+} or by an inhomogeneous potential distribution.

417

418 Despite that the combination of the hypotheses of electroneutrality and (induced)
419 turbulence can explain the observed observations, we stress the tentative character of all
420 these speculations.

421

422 4.3 Practical strategies to overcome the anomalous convex behaviour

423 We aim now at providing an estimation of the conditions under which acb can be a
424 problem for AGNES when the intensity current is taken as the response function. It is
425 clear from what has been commented that the acb appears for high [Zn²⁺] and low ionic
426 strength. Indeed, we have seen how for ionic strength 0.05M, acb appears around
427 [Zn⁰] $\approx 9 \times 10^{-4}$ M, for a ionic strength 0.10M the acb started at [Zn⁰] ≈ 0.02 M and for a
428 ionic strength 0.5M acb is seen at [Zn⁰] ≈ 0.05 M. However, for conditions before
429 sluggish stripping, the acb can be remediated with longer t_1 (observation 5). As seen in
430 Fig 2, for the same value of [Zn⁰]=Y[Zn²⁺], we can have experiments with acb (such as
431 the four upper curves) and experiments without acb (which act as reference, see thick
432 line) just by changing the deposition time. Fig 10 classifies a large number of
433 experiments (having reached the target condition of Nernstian equilibrium and absence
434 of gradients by the end of the first stage) according to the observed acb through a visual
435 inspection similar to that already described for Fig 2. Indeed, the procedure of plotting
436 the normalised current (or $\eta(t_2)$) can serve as diagnosis of the acb. If acb is found,
437 taking the faradaic current at a fixed t_2 as response function of AGNES is no longer
438 valid.

439

440 Consistent with observations 1 and 5, we see in Fig 10 that acb predominates for higher
441 [Zn⁰] and shorter times, while the standard behaviour for the stripping is seen for lower
442 [Zn⁰] and longer deposition times. So, from a practical point of view, if a given

443 application of AGNES to a sample exhibits acb (diagnosed from the divergence of the
444 sample measurement from the reference curve which can be obtained –for instance- in
445 the calibration at not very high [Zn⁰]), one can choose to reduce Y , to increase t_1 or a
446 combination of both.

447

448 Taking into account the results reported in this work, we detail below 4 strategies to
449 apply AGNES in systems with a large Zn concentration and/or low ionic strength (i.e.
450 when acb can appear or has been diagnosed):

451

452 a) Reduction of the prescribed preconcentration factor:

453 The reduction of Y is very easy to implement and very convenient, because the ensuing
454 required deposition time (t_1) is proportionally reduced. Indeed, for gains below $Y=5$ and
455 drop 1 in the HMDE, a deposition time of 50 s is more than sufficient –if other
456 complications do not operate- to obtain the equilibrium goal. The only caution is that
457 the response current has to be sufficiently higher than the blank to ensure a
458 determination above the limit of quantification. We have previously reported the need
459 of reducing the Y along metal titrations experiments of humic acid, when covering a
460 wide range of metal concentration [8].

461

462 b) the use of the charge as response function:

463 Alternatively, in order to avoid acb, we can exploit observation 2. For those systems not
464 exhibiting adsorption, the value of the stripped charge can be used as response function
465 (instead of the faradaic current) as long as other phenomena of poor reproducibility do
466 not appear. As seen in fig 11, Q is linear with the product of $[Zn^0]=Y [Zn^{2+}]$ for a range
467 of experiments exhibiting or not acb (experiments from Fig 10) –being the slope, of

468 $2.05 \times 10^{-3} \text{ C M}^{-1}$, the value of the experimental η_Q . Thus, once more, we see that the
469 charge can be useful even when the intensity current fails. This supports the practical
470 possibility of using AGNES by taking Q as the response function: determining η_Q from
471 calibration plots like Fig 11, measuring the charge in the sample and computing the free
472 concentration with eqn (8). Shorter interval times (Δt) will lead to more accurate
473 determinations of the charge.

474

475 c) the use of longer deposition times:

476 As commented above in subsection a), and in accordance with observation 5, there is a
477 range of conditions (given by the analyte concentration, the desired gain and the ionic
478 strength) where the increasing of the deposition time restores the linearity of the current
479 response with Zn^{2+} concentration.

480 From diagram of Fig 10 we can formulate a rough rule of thumb, for ionic strength 0.05
481 M, of t_1 having to be

$$482 \quad t_1 > 2 \times 10^5 \times Y \times [\text{Zn}^{2+}] \quad (12)$$

483 in order to avoid acb. Of course, the application of this strategy is rather limited because
484 the resulting “recommended” deposition time from the previous equation might be
485 prohibitively long.

486

487 d) use higher ionic strength:

488 Increasing the ionic strength decreases the acb effect, thus enlarging the range of
489 linearity for the current. So, when working with some synthetic solutions and large Zn
490 concentrations, it might be easier to work with increased concentrations of background
491 electrolyte. However, this strategy has also a limited use, given that for many samples,

492 the changing of the ionic strength can be critical (e.g. when measuring Zn in river water,
493 the alteration of the ionic strength changes the speciation).

494

495 **5. Conclusions**

496 The standard application of AGNES to determine high Zn concentrations can be
497 hindered by the anomalous convex behaviour (acb). This problem is more critical at
498 lower ionic strengths, given that there is a trend towards the standard behaviour at
499 sufficiently large concentrations of background electrolyte. The anomalous behaviour
500 can be diagnosed by plotting the normalised current $\eta(t_2)$ along the stripping stage and
501 comparison with a reference curve. We have described (section 4.1) a set of phenomena
502 associated to this acb, such as the appearing of initial enhanced currents (with respect to
503 the diffusion limited reference) or the sluggish stripping.

504

505 We speculate (see section 4.2) that the low conductivity in solution (which is also a
506 levelling-off mechanism) and an induced turbulence can explain the previous
507 observations regarding acb, but further work should elucidate this issue.

508

509 AGNES can be applied to determine high Zn concentrations, provided that care is taken
510 to diagnose or circumvent the acb of the stripping currents. As strategies to avoid the
511 loss of linearity in the current for high ratios of analyte concentration over background
512 electrolyte concentration, we suggest (see section 4.3): a) the reduction of the applied
513 gain; b) the use of the charge as response function (see also section 2); c) the use of
514 longer deposition times (even if this is a limited measure) and d) use of higher ionic
515 strength (also a limited measure).

516

517 **Acknowledgments**

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520 CTM2006-13583) and from the “Comissionat per a Universitats i Recerca del
521 Departament d’Innovació, Universitats i Empresa de la Generalitat de Catalunya” and
522 from the European Social Fund.

523

524

525 **6. Appendix**

526

527 We discuss here with more detail some of the hypotheses listed in section 4.2 which
528 could participate in explaining the observations reported in section 4.1.

529

530 iii) The irreversibility hypothesis. There is a wealth of literature indicating that the
531 system $Zn^0 \rightarrow Zn(II)$ is irreversible. One objection to this being the cause of the
532 observations 1 to 6 is that the applied reoxidation potential (E_2) is much more positive
533 than E^0 and irreversibility should decrease with the applied overpotential. Moreover, a
534 very similar sluggish stripping has been observed for Cd (see Fig), which is an element
535 considered to behave reversibly.

536

537 iv) The “island” hypothesis. Some authors [12,14] have considered the temporary
538 formation of a “concentration inhomogeneity”. One might think that -during the first
539 moments of the deposition step- large amounts of Zn^0 gather inside the amalgam close
540 to the surface and form a kind of island in another phase or “association state”. For
541 instance, BenBassat et al. [13] interpreted the extra peaks in Linear Anodic Stripping
542 Voltammetry (L-ASV) found with Cd as due to the reoxidation of such islands. We
543 have repeated a similar experiment to that of BenBassat *et al.*, with Cd and Zn and with

544 a deposition stage as that of AGNES (i.e. at a moderate Y). Our AGNES-like L-ASV
545 voltammograms (Fig 8 and) for short deposition times exhibit what could be called an
546 extra peak. Thus, our results in these figures could be consistent with this hypothesis by
547 explaining that when t_1 is very long, the island -which is metastable- disappears[15]. In
548 this way, we could understand observation 3 (sluggish stripping) as a kind of kinetic
549 limitation of the re-dissolution of the island (analogous to that seen in fig. 8 of ref. [31]),
550 and we could explain observation 4 (enhanced currents) as part of the deposited Zn⁰
551 being closer to the surface, but some other phenomenon must be involved, given that
552 the existence of the islands can be invoked both to explain the slowing down and the
553 speeding up of the standard stripping. Be as it may, observation 2 (Q being correct for
554 concentrations when the enhanced currents occur) indicates that AGNES conditions
555 have been attained, so we need to accept that the desired Zn⁰ concentration has been
556 attained uniformly inside the drop at the fixed value $Y[\text{Zn}^{2+}]$. If there were other phases
557 or compounds, we would record larger charges for longer t_1 (and we have seen this is
558 NOT the case in many experiments; e.g. in 2 or 5). To render the island hypothesis
559 compatible with the experimentally found charge corresponding to AGNES
560 requirements, we would need to accept that some Zn⁰ is converted into a form which
561 counts for the total activity inside the drop, but with very different availability for
562 different experiments.

563

564 v) The shell hypothesis. The optical observation of a shell around the mercury drop,
565 described in observation 7, suggests the possibility of existence of similar shells for
566 lower concentrations (despite we have not been able to see them with our microscope).
567 The sluggish stripping, observation 3, could be explained by assuming that a shell of
568 Zn⁰ and/or other materials gathers -without mechanical contact with Hg, so that

569 observation 2 holds- around the drop hindering the stripping. First objection: why does
570 this hypothetical shell not hinder the deposition stage?. Another objection: the
571 enhancement seen in observation 4 cannot be explained by this hypothesis. A third
572 objection to this hypothesis is that this shell effect should imply the usually large initial
573 reoxidation current (collapsing with the reference curve) at very short stripping times
574 (because diffusion in the thin solution layer in between the drop and the shell would not
575 be initially hindered) which has not been observed .

576
577 vi) Intermetallic formation. The formation of intermetallic compounds of Zn and Cu
578 could produce a decreased AGNES intensity current (sluggish stripping), given that it is
579 well known that more positive potentials than those actually applied (close to that of
580 Cu⁰) are required for their reoxidation [32]. From the maximum declared content of Cu⁰
581 in the used polarographic mercury (5×10⁻⁶%), one estimates [Cu⁰]_{from amalgam} =1×10⁻⁵
582 M. To this amount one should add the Cu⁰ deposited from the solution which will
583 depend on the used [KNO₃] and deposition time the KNO₃. From the maximum content
584 of Cu in KNO₃ declared by the supplier (5×10⁻⁷%), we compute a typical [Cu²⁺]_{≈10⁻⁹}
585 M, from which the deposited concentration can be estimated as

$$586 \quad [Cu^0]_{\text{from solution}} = \frac{D_{Cu^{2+}} [Cu^{2+}] 4\pi r_0^2}{\delta^4 / 3 \pi r_0^3} t_1 \approx 5 \times 10^{-10} \text{ M} \quad (13)$$

587 where δ is the diffusion layer thickness (around 2×10^{-5} m) and t_1 is taken around 1000 s.
588 We conclude that the amount of Cu⁰ coming from the solution is negligible in front of
589 the one in the amalgam, but, in turn, [Cu⁰]_{from amalgam} is negligible in front of the [Zn⁰]
590 concentrations used when the linearity of AGNES response is lost. Given that in our
591 experiments [Zn⁰]_{>>}[Cu⁰]_{from amalgam}, the most likely intermetallic compounds to
592 interfere in AGNES measurements are CuZn or CuZn₃ (the most unfavourable
593 case)[33]. As AGNES intensity current (or charge) depends on [Zn⁰] and its variation,

594 due to the possible formation of intermetallic compounds, is negligible
595 ($3 \times [\text{Cu}^0] \ll [\text{Zn}^0]$), we conclude that –for these cases- their formation cannot explain the
596 sluggish stripping. We can even argue more directly noting that –if Cu^0 coming from
597 solution is negligible-, once all Cu^0 (from amalgam contamination) becomes the
598 intermetallic CuZn_3 , the accumulation of Zn^0 will proceed along the first stage of
599 AGNES until $[\text{Zn}^0]$ will reach the desired $Y[\text{Zn}^{2+}]$ and –given that CuZn_3 is not
600 reoxidized at the used stripping potentials for Zn^0 - the response function will be the
601 same as if Cu was not present. Work is in progress in our laboratory to study the impact
602 of intermetallic compounds on AGNES.

603

604

605

606

Table

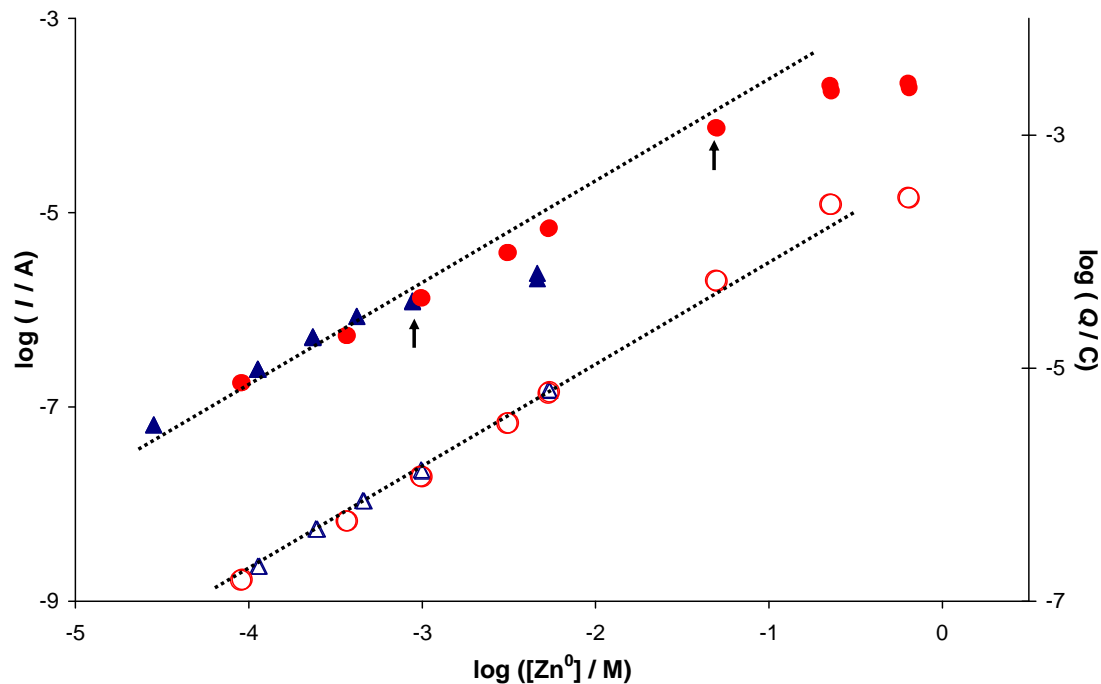
Table 1: Simplest potential program of AGNES.

potential	deposition at E_1 for prescribed gain (Y)	deposition at E_1 (rest period)	E_2 in stripping or second stage (gain Y_2)
time	$t_1 - t_w$	t_w	t_2
stirring	on	off	off

607

608

609 **Figures**

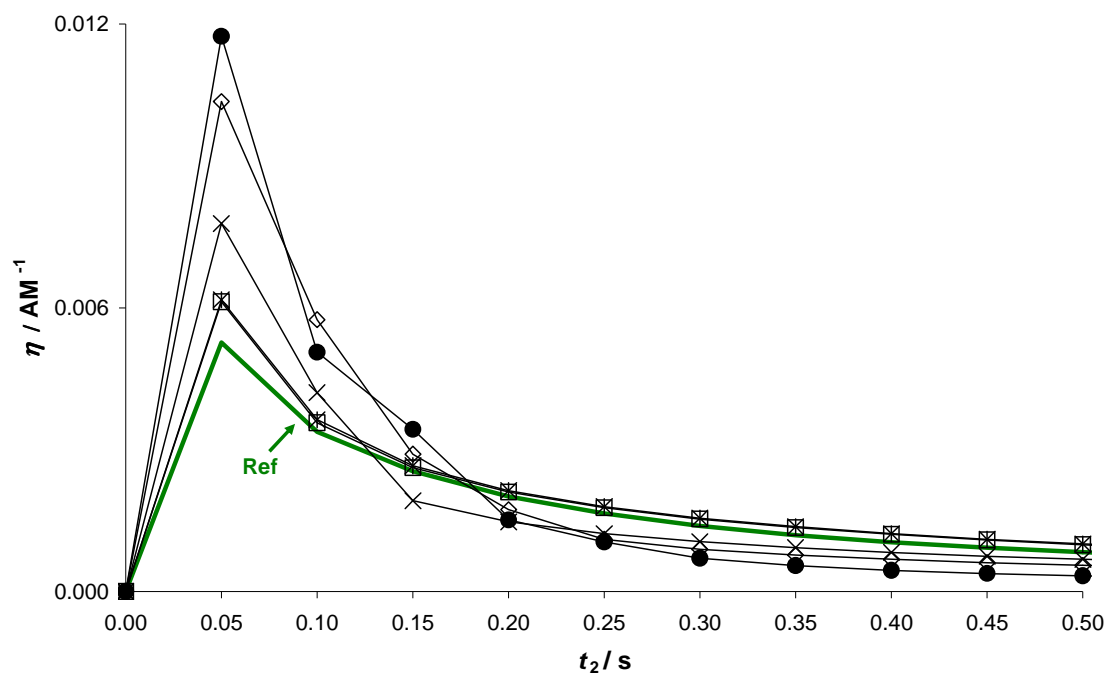


610

611

612 Fig 1: Calibration plot of the logarithm of the faradaic current (full markers referred to
613 the left axis) and of the logarithm of the charge (open markers referred to the right axis)
614 vs. Zn concentration inside the amalgam. I is measured at $t_2=200$ ms, while Q is
615 computed with $\Delta t= 50$ ms in eqn. (9). Makers: triangle for $[KNO_3]=0.05$ M and circle
616 for $[KNO_3]=0.5$ M. Settings: $Y= 50$; $t_1= 400$ s ; $t_w=50$ s. The arrows indicate the first
617 concentration of each $[KNO_3]$ which clearly exhibits acb (via an analysis as that
618 performed in Fig 2) and also corresponds to the loss of linearity of current in this
619 calibration plots. The dotted straight lines are lines of slope 1 as expected from eqn. (4).

620



621

622 Fig 2: Plot of the normalised current $\eta(t_2) = \frac{I(t_2)}{Y[Zn^{2+}]}$ -see eqn. (3)- showing how the

623 anomalous convex behaviour in the stripping current decreases with increasing

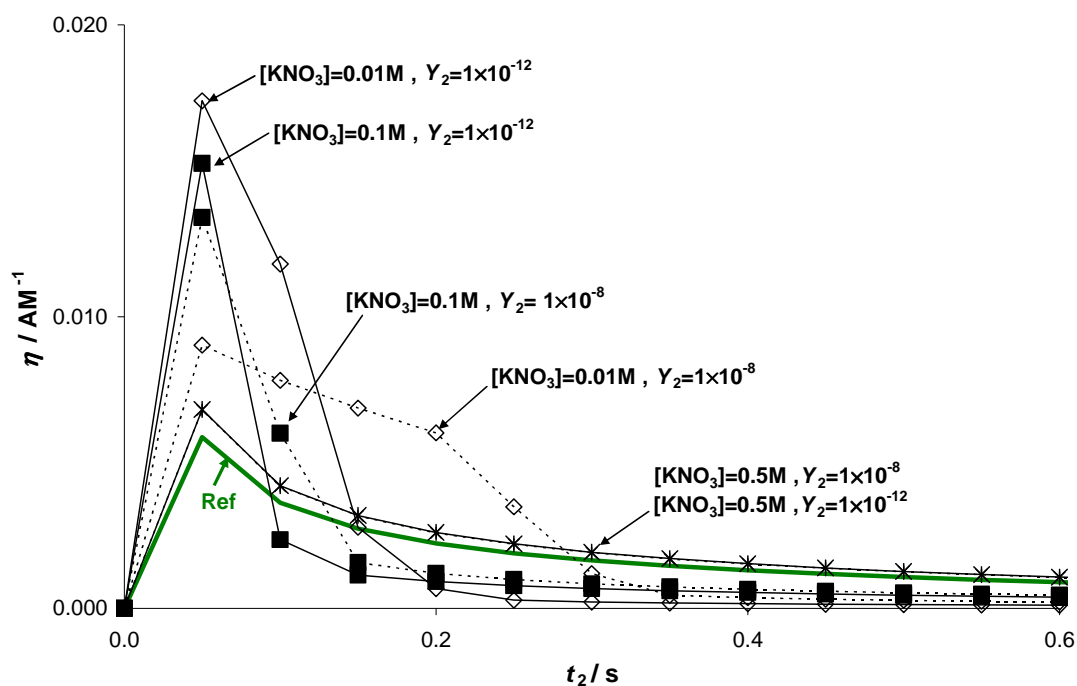
624 deposition time. Markers for experimental data: (●) $t_1=400s$, (◇) $t_1=800s$, (×) $t_1=1200s$,

625 (□) $t_1=2000s$ and (*) $t_1=2200s$. Other conditions: $[Zn^{2+}]=1.3 \times 10^{-5} M$; $Y=50$; $[Zn^0]=$

626 $6.5 \times 10^{-4} M$, $[KNO_3]=0.1M$. The thick line indicates the “reference” (or typical)

627 stripping decay, which has been achieved with $[Zn^{2+}] = 3.0 \times 10^{-6} M$, $Y=50$.

628

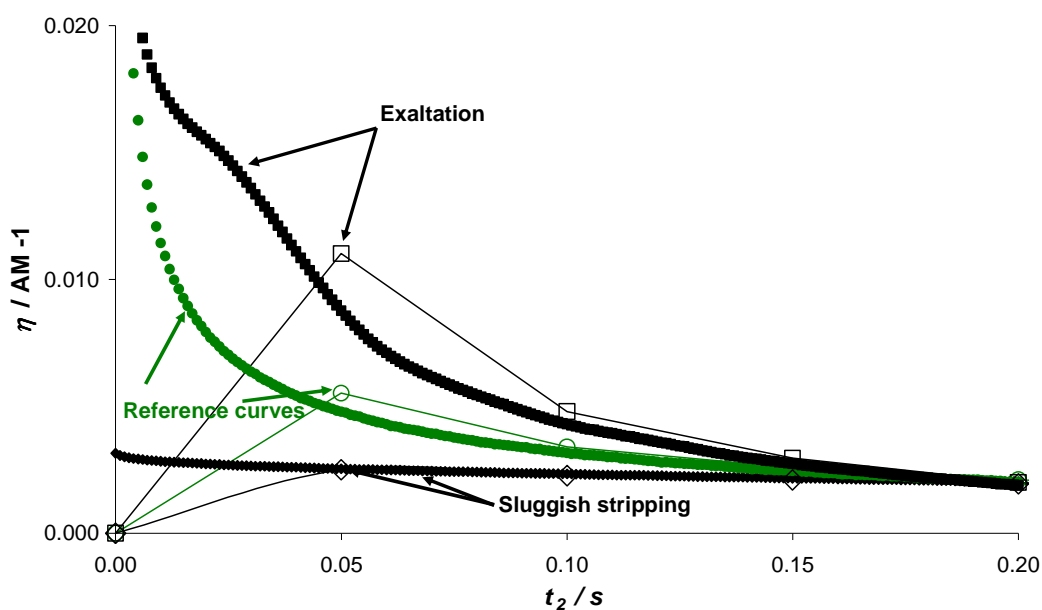


629

630 Fig 3: Plot of the normalised currents vs. t_2 for $[Zn^{2+}] = 5.8 \times 10^{-5} M$, different stripping
 631 potentials E_2 and $[KNO_3]$. Markers for experimental data: (—) $E_2 = -0.5956V, -0.6025V$
 632 and $-0.6102V$ for $[KNO_3] = 0.01M, 0.1M$ and $0.5M$ respectively, corresponding to
 633 $Y_2 = 1 \times 10^{-12}$; (---) $E_2 = -0.7139V, -0.7208V, -0.7285V$ for $[KNO_3] = 0.01M, 0.1M$ and $0.5M$
 634 respectively, corresponding to $Y_2 = 1 \times 10^{-8}$; (\diamond) $[KNO_3] = 0.01M$; (\blacksquare) $[KNO_3] = 0.1M$; ($*$)
 635 $[KNO_3] = 0.5M$. Settings: $t_1 = 400s, t_w = 50s, Y = 50, [Zn^0] = 2.9 \times 10^{-3} M$. The thick line
 636 indicates the reference stripping decay obtained with $Y = 1$.

637

638

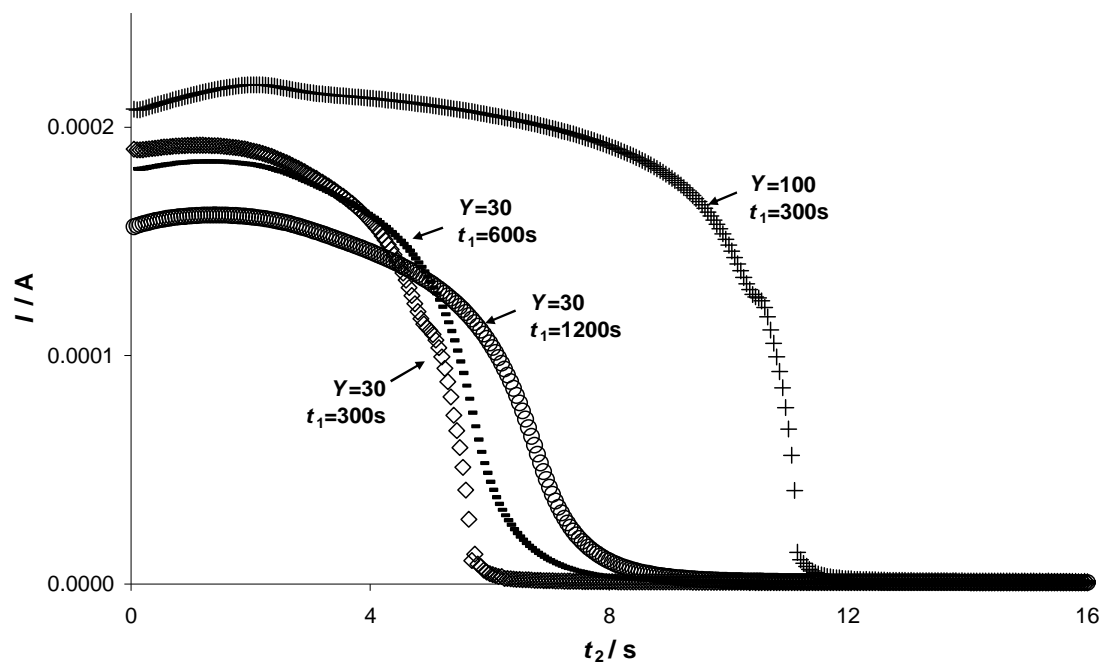


639

640 Fig 4: Plot of the normalized current η vs. t_2 showing the two types of acb: exaltation
641 (square marker) and sluggish stripping (diamond marker) compared with a reference
642 curve (circle marker). Open markers indicate sample time $\Delta t = 50$ ms and full markers
643 indicate $\Delta t = 1$ ms. Parameters: (o) $Y=1$, $t_l = 50$ s, $\Delta t = 50$ ms; (●) $Y=1$, $t_l = 50$ s, $\Delta t = 1$ ms;
644 (□) $Y=2$, $t_l = 50$ s, $\Delta t = 50$ ms; (■) $Y=2$, $t_l = 50$ s, $\Delta t = 1$ ms; (◇) $Y=50$, $t_l = 400$ s, $\Delta t = 50$ ms;
645 (◆) $Y=50$, $t_l = 400$ s, $\Delta t = 1$ ms. Settings: $[\text{Zn}^{2+}] = 1.90 \times 10^{-4} \text{M}$, $[\text{KNO}_3] = 0.01 \text{M}$, $t_w = 50$ s.

646

647



648

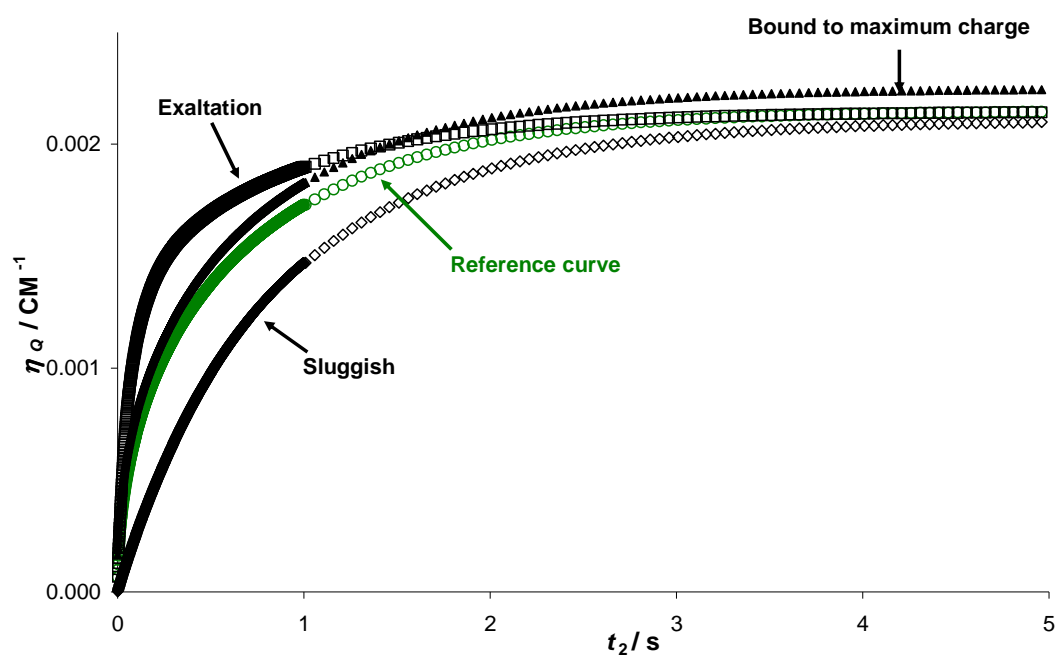
649 Fig 5: Plot showing the sluggish stripping currents for huge concentrations. $Y=30$

650 (several deposition times: (\diamond), $t_1=300s$, (—) $t_1=600s$ (o) and $t_1=1200s$) and $Y=100$ ((+)

651 $t_1=300s$). Settings: $[Zn^{2+}] = 1.36 \times 10^{-2} M$; $[KNO_3] = 0.05 M$.

652

653



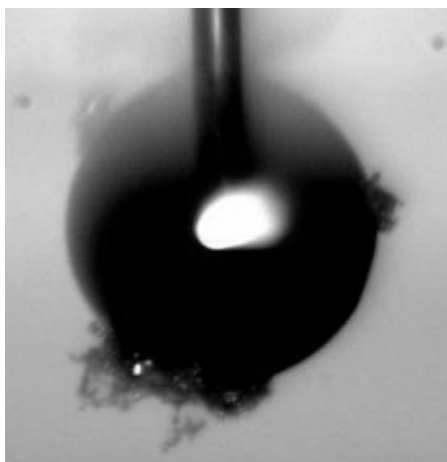
654

655 Fig 6: Plot of the normalized charge η_Q (see eqn. (7)) vs. t_2 for the data of Fig 4 (same
656 markers) with $\Delta t=1$ ms for the first second of the stripping. The curve that is always
657 below the reference curve is classified as sluggish stripping while the curve that –at
658 some point- overcomes the reference one is classified as corresponding to enhanced
659 currents. The exaltation even overcomes the solid triangle series that stands for the
660 maximum expected charge (under diffusion limited conditions) computed applying eqn.
661 (11) to the data of the reference experiment.

662

663

664 a)



b)



665

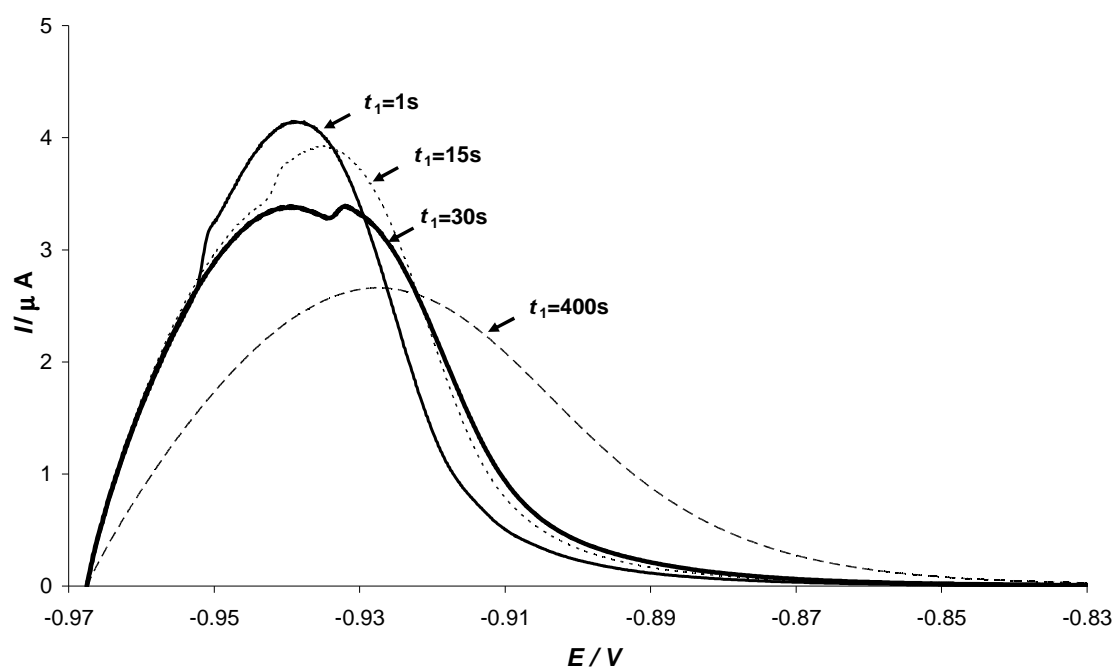
666

667 Fig 7: Photographs of the mercury drop electrode obtained in an AGNES experiments
668 with diffusion limited deposition potential ($Y = 10^{10}$) at $[\text{KNO}_3] = 0.05 \text{ M}$. a) $[\text{Zn}^{2+}] =$
669 $3.76 \times 10^{-2} \text{ M}$ and $t_1 = 400 \text{ s}$. b) $[\text{Zn}^{2+}] = 2 \times 10^{-2} \text{ M}$ and $t_1 = 1800 \text{ s}$

670

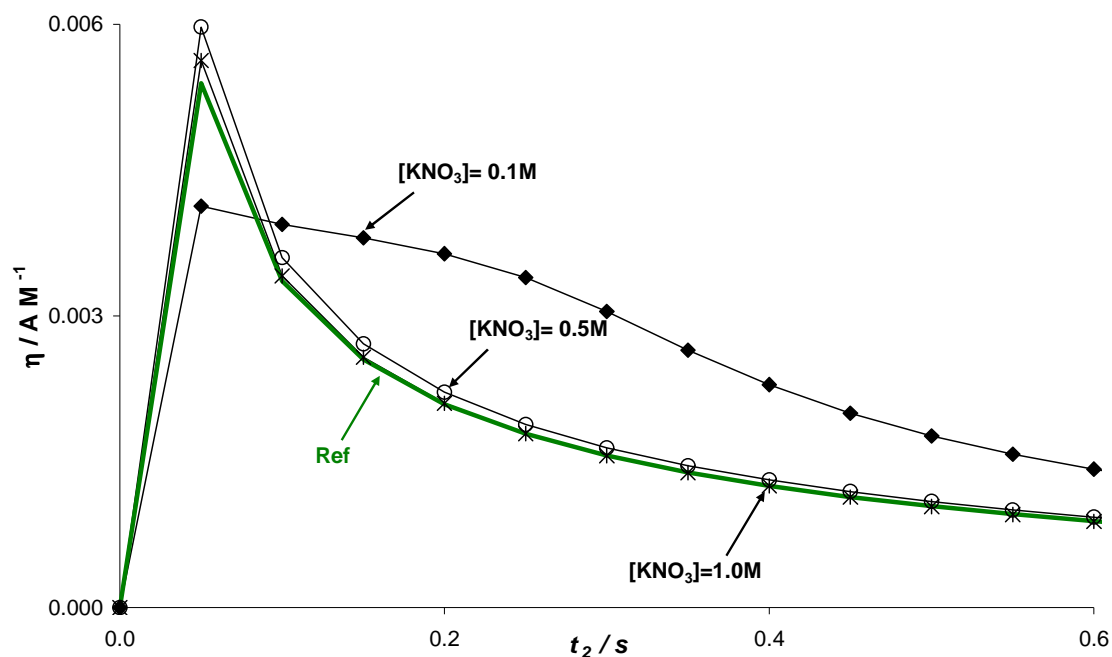
671

672



673

674 Fig 8: Linear Anodic Stripping Voltammogram (L-ASV) of a solution with
675 [Zn²⁺]=4.81×10⁻³M and several deposition times. Thin continuous line t₁= 1s, thin
676 dotted line t₁= 15s, thick continuous line t₁= 30s, dashed line t₁= 400s. Deposition stage
677 with Y=1. Linear stripping with scan rate: 17 mV/s and ΔE= 0.15 mV, [KNO₃]= 0.05M.



678

679 Fig 9: Plot of η vs. t_2 for Cadmium with different ionic strengths . Parameters: $[\text{Cd}^{2+}] =$

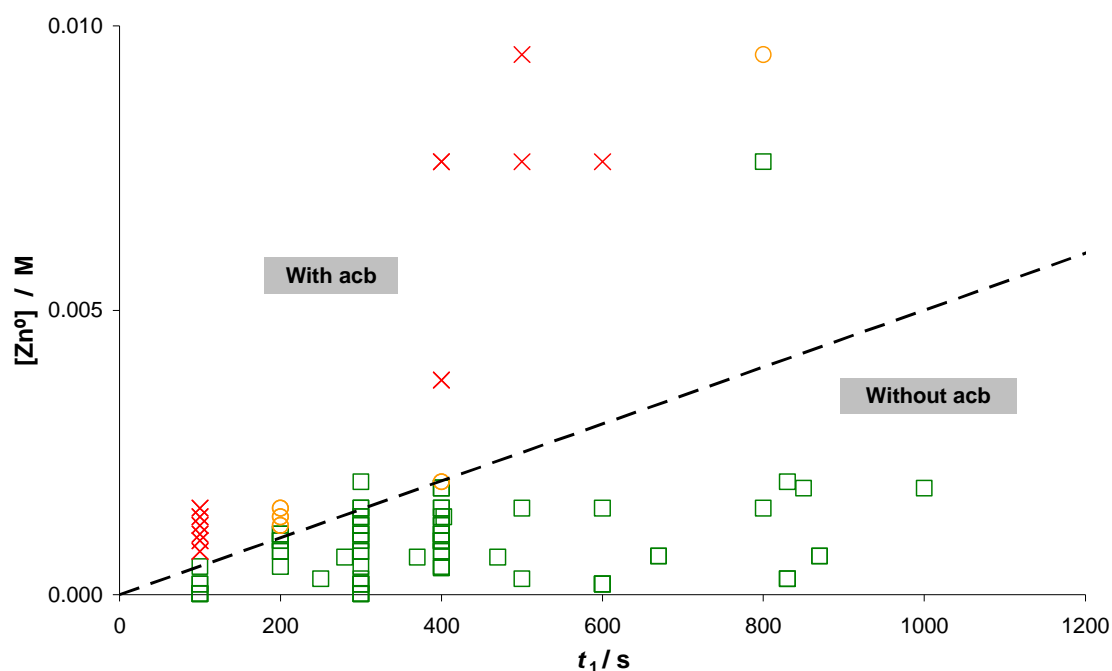
680 $3.2 \times 10^{-3} \text{M}$, $Y=12$, (\blacklozenge) $[\text{KNO}_3]=0.1\text{M}$ and $t_1=400\text{s}$; (\circ) $[\text{KNO}_3]=0.5\text{M}$ and $t_1=1200\text{s}$; ($*$)

681 $[\text{KNO}_3]=1.0\text{M}$ and $t_1=400\text{s}$. Reference curve: $[\text{Cd}^{2+}] = 3.2 \times 10^{-3} \text{M}$, $Y=8$, $t_1=400\text{s}$.

682

683

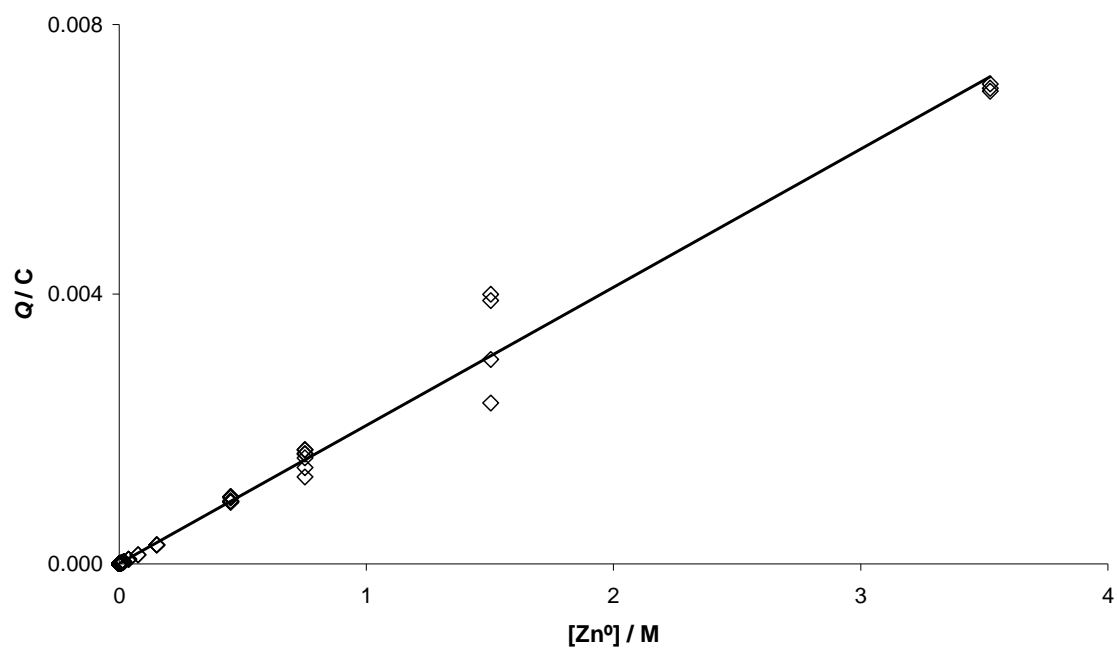
684



685

686 Fig 10: Nomograph indicating the appearance of the acb in the stripping stage for
687 different combinations of the product $Y [Zn^{2+}]$ (i.e. aimed $[Zn^0]$) and deposition times
688 (t_1) in the simplest potential program of AGNES for more than 50 experiments with
689 various Zn^{2+} concentrations, $[KNO_3]= 0.05M$. Markers: (□) curve indicates experiments
690 where acb effect was not detected; (×) indicates experiments where acb could be clearly
691 identified; (○) indicate experiments where the acb effect (if existent) was very mild.

692

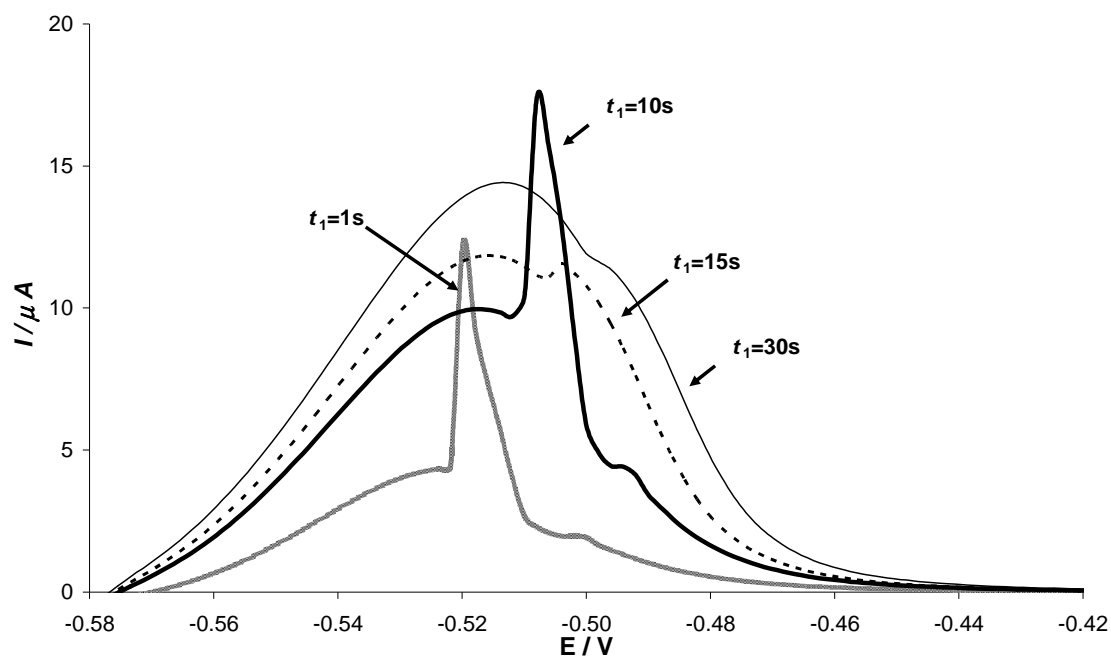


693

694 Fig 11: Plot of the total stripped charge ($\Delta t = 50$ ms in eqn. (9)) vs. $[Zn^0] = Y[Zn^{2+}]$ for the
695 same experiments analyzed in fig 10 and a few more at higher concentrations. The slope
696 of the line (2.05×10^{-3} C/M) is the factor η_Q which can be used together with the
697 measured charge when the acb affects the faradaic current.

698

699



700

701 Fig 12: Linear Anodic Stripping Voltammogram (L-ASV) of a solution with [Cd²⁺]=
702 1.34×10⁻³M. Deposition potential corresponding to Y=20. Deposition times: (---) t₁=
703 1s; (—) t₁= 10s; (- - -) t₁= 15s; (-) t₁= 30s. Linear stripping with scan rate: 0.017V/s and
704 ΔE= 0.00015V. [KNO₃]= 0.05M.

705

706 References

707

708

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